

The geometry of Haldane limits

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Nonlinear σ -models

- Entered physics via low-energy QCD with the work of M.Gell-Mann and M.Levy (1960)
- Describe the scattering of Goldstone bosons in 4D, for example π -mesons in the case $\frac{SU(2) \times SU(2)}{SU(2)}$ (u, d quarks)
- Renormalizable in 2D

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Therefore provides a method for the exploration of target-space geometry
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- QFT requires regularization preserving the symmetries
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- F.D.M.Haldane, 1983

$SU(2)$, representation of spin $\mathbf{s} = \frac{m}{2}$ on $\text{Sym}(V_{\text{fund}}^{\otimes m})$ with Hamiltonian

$$\mathcal{H} = \sum_i \vec{S}_i \cdot \vec{S}_{i+1}$$

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The goal / result

- The goal is to construct the spin chain with target space $\frac{U(N)}{U(n_1) \times \dots \times U(n_m)}$.

The Hamiltonian

$$H = \sum_{i=1}^L \sum_{k=1}^{m-1} d_k \vec{S}_i \cdot \vec{S}_{i+k}$$

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where $z \in \mathbb{C}P^{N-1}$ and

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Elements of symplectic geometry

- Phase space \mathcal{N} is a symplectic manifold.
- ω is a non-degenerate closed 2-form: $d\omega = 0$
- $G \curvearrowright \mathcal{N}$, the moment map $\mu : \mathcal{N} \rightarrow \mathfrak{g}$
- Equivariance:
$$\mu(g \circ x) = \text{Ad}_g \mu(x) \equiv g\mu(x)g^{-1}$$
- Hamiltonians: $d\mu_a = i_{X_a}\omega, \quad a \in \mathfrak{g}$
- Simplest example: angular momentum

$$\mathcal{N} = \mathbb{R}^6, \quad G = SO(3), \quad \omega = d\vec{r} \wedge d\vec{p}, \quad \vec{L} = \vec{r} \times \vec{p}.$$

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The A.Borel-A.Weil-R.Bott theorem

- Let V be a representation of $U(N)$ with highest weight $\vec{\lambda}$.
- It can be built on the space of sections of a holomorphic fiber bundle

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- In special cases the base can be reduced:



$$\frac{U(4)}{U(3) \times U(1)} = \mathbb{C}P^3$$



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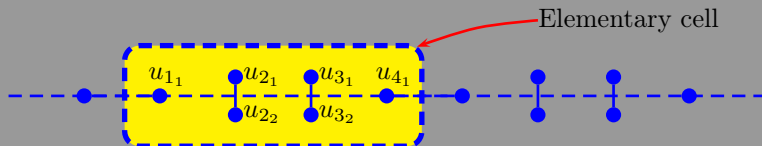


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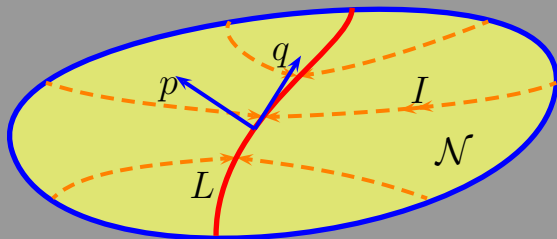
A picture of the spin chain



The spin chain for the coset $\frac{U(6)}{U(2) \times U(2) \times U(1) \times U(1)}$

The general setup

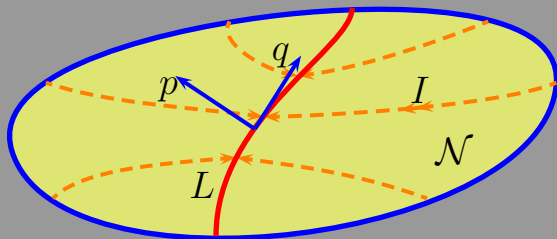
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- Consider a function I which has a minimum on a Lagrangian submanifold $L \subset \mathcal{N}$.
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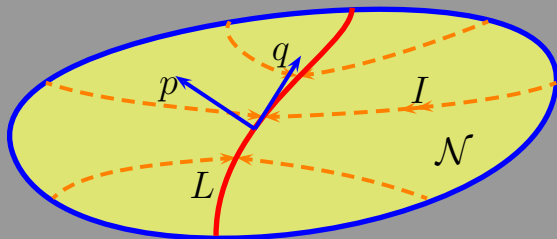
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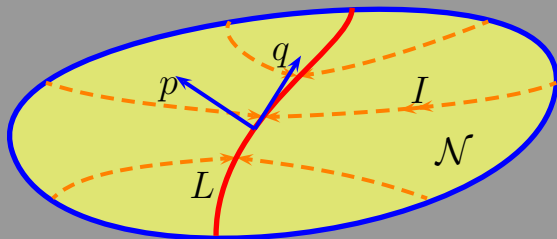
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The metric

- The metric in the normal directions to L : the Hessian $h_{ij} = \frac{\partial^2 I}{\partial x^i \partial x^j}$

The metric on L

$$g_{ij} = \omega_{im} \cdot \left[\left(\frac{\partial^2 I}{\partial x^2} \right)^{-1} \right]^{mn} \cdot \omega_{nj} = \omega_{im} h^{mn} \omega_{nj}$$

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The θ -term

- Construct the Lagrangian embedding

$$i : \mathcal{F}_{n_1, \dots, n_m} \hookrightarrow G_{n_1} \times \dots \times G_{n_m}$$

- Ingredients for the θ -term:

$$r_k = i^*(c_1(\mathcal{O}_{G_{n_k}}(1))), \quad \sum_{k=1}^m r_k = 0.$$

The θ -term

$$\Omega = \frac{1}{m} \left(\sum_{k=1}^m k \cdot r_k \right)$$

- Hence $\theta = \frac{2\pi}{m}$. Permuting the sites of the spin chain changes the θ -term in $\mathbb{H}^2(\mathcal{M}, \mathbb{Z}_m)$!

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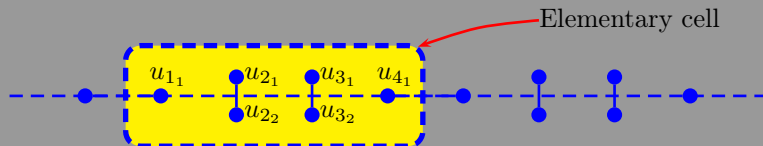
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- Relation to the action of Weyl group on Schubert cells / Bruhat decomposition

The picture revisited



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Questions / Answers

- Constructed a spin chain for a σ -model with target space a flag manifold.
- Universal expressions for the metric and θ -term.
- Is there a **mod m** periodicity of the mass gap?
- Is there an efficient way to describe the σ -models numerically using spin chains?

Questions / Answers

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